

Vibration Mechanics Hw #6

(Multiple-Degree-of-Freedom Systems)

145 Pts

Issued: May 19 (FRI), 2023

Due: June 05 (MON), 2023

1. Rao P. 6.2 (please refer to Figure 6.19) 15 Pts

Equations of motion:

$$J_0 \ddot{\theta} = -2k \left(\frac{\ell}{4} \theta - x_1 \right) \frac{\ell}{4} - c \left(\frac{\ell}{4} \dot{\theta} - \dot{x}_1 \right) \frac{\ell}{4} - 3k(\theta \ell) \ell + M_t$$

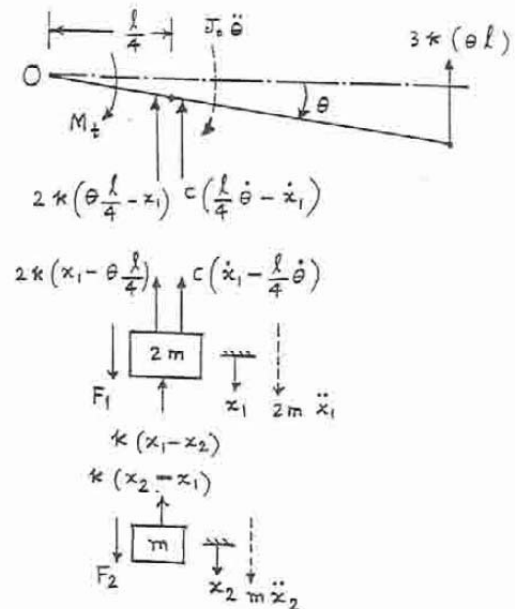
$$2m \ddot{x}_1 = -2k \left(x_1 - \frac{\ell}{4} \theta \right) - c \left(\dot{x}_1 - \frac{\ell}{4} \dot{\theta} \right) - k(x_1 - x_2) + F_1$$

$$m \ddot{x}_2 = -k(x_2 - x_1) + F_2$$

where $J_0 = \frac{1}{3} (2m) \ell^2 = \frac{2}{3} m \ell^2$

These equations can be stated in matrix form as:

$$\begin{bmatrix} \frac{2}{3} m \ell^2 & 0 & 0 \\ 0 & 2m & 0 \\ 0 & 0 & m \end{bmatrix} \begin{Bmatrix} \ddot{\theta} \\ \ddot{x}_1 \\ \ddot{x}_2 \end{Bmatrix} + \begin{bmatrix} \frac{c \ell^2}{16} & -\frac{c \ell}{4} & 0 \\ -\frac{c \ell}{4} & c & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{Bmatrix} \dot{\theta} \\ \dot{x}_1 \\ \dot{x}_2 \end{Bmatrix} + \begin{bmatrix} \frac{25k \ell^2}{8} & -\frac{k \ell}{2} & 0 \\ -\frac{k \ell}{2} & 3k & -k \\ 0 & -k & k \end{bmatrix} \begin{Bmatrix} \theta \\ x_1 \\ x_2 \end{Bmatrix} = \begin{Bmatrix} M_t \\ F_1 \\ F_2 \end{Bmatrix}$$



From strength of materials, the deflection of the cantilever beam shown is given by

$$w(x)|_{in AB} = \frac{Fx^2}{6EI}(-x + 3a) \quad \dots (E_1)$$

$$w(x)|_{in BC} = \frac{Fa^2}{6EI}(-a + 3x) \quad \dots (E_2)$$

Apply $F_1 = 1, F_2 = F_3 = 0$: $a_{11} = (F=1, x=l, a=l \text{ in } (E_1)) = l^3/(3EI)$

$$a_{21} = (F=1, x=2l, a=l \text{ in } (E_2)) = 5l^3/(6EI)$$

$$a_{31} = (F=1, x=3l, a=l \text{ in } (E_2)) = 4l^3/(3EI)$$

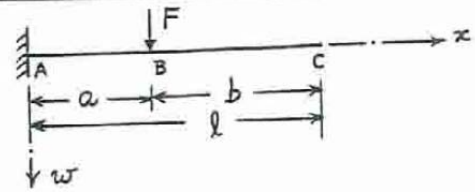
Similarly apply $F_2 = 1, F_1 = F_3 = 0$ to get a_{22}, a_{32}, a_{12} and $F_3 = 1, F_1 = F_2 = 0$ to get a_{33}, a_{13}, a_{23} . Result is

$$[a] = \frac{l^3}{EI} \begin{bmatrix} 1/3 & 5/6 & 4/3 \\ 5/6 & 8/3 & 14/3 \\ 4/3 & 14/3 & 9 \end{bmatrix}$$

Equations of motion:

$$[m] \ddot{\vec{w}} + [k] \vec{w} = \vec{0} \quad \text{or} \quad [a][m] \ddot{\vec{w}} + \vec{w} = \vec{0}$$

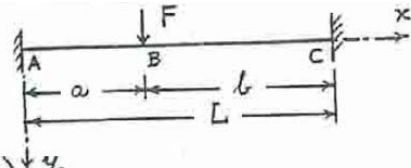
with $[m] = \begin{bmatrix} m_1 & 0 & 0 \\ 0 & m_2 & 0 \\ 0 & 0 & m_3 \end{bmatrix}$ and $\vec{w} = \begin{Bmatrix} w_1 \\ w_2 \\ w_3 \end{Bmatrix}$.



Deflection of a fixed-fixed beam is

$$y(x) \Big|_{in AB} = \frac{Fb^2x^2}{6EI L^3} \{-x(3a+b) + 3aL\} \dots (E_1)$$

$$y(x) \Big|_{in BC} = \frac{Fa^2(L-x)^2}{6EI L^3} \{-(L-x)(3b+a) + 3bL\} \dots (E_2)$$



Apply $F_1 = 1, F_2 = F_3 = 0$:

$$a_{11} = (F=1, a=l, b=3l, x=l, L=4l \text{ in } (E_1)) = 9l^3/(64EI)$$

$$a_{21} = (F=1, a=l, b=3l, x=2l, L=4l \text{ in } (E_2)) = l^3/(6EI)$$

$$a_{31} = (F=1, a=l, b=3l, x=3l, L=4l \text{ in } (E_2)) = 13l^3/(192EI)$$

Similarly apply $F_2 = 1, F_1 = F_3 = 0$ to get a_{22}, a_{32}, a_{12} and

$F_3 = 1, F_1 = F_2 = 0$ to get a_{33}, a_{13}, a_{23} . Result is

$$[a] = \frac{l^3}{EI} \begin{bmatrix} 9/64 & 1/6 & 13/192 \\ 1/6 & 1/3 & 1/6 \\ 13/192 & 1/6 & 9/64 \end{bmatrix}$$

Coordinates of the bob are $(x + l \cos \theta, l \sin \theta)$

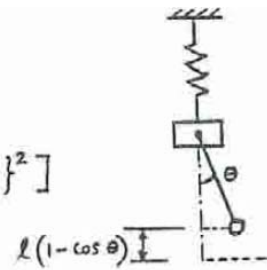
$T = \text{kinetic energy} = \text{kinetic energy of slider}$
 $+ \text{kinetic energy of bob}$

$$= \frac{1}{2} m \left(\frac{dx}{dt} \right)^2 + \frac{1}{2} m \left[\left\{ \frac{d}{dt} (x + l \cos \theta) \right\}^2 + \left\{ \frac{d}{dt} (l \sin \theta) \right\}^2 \right]$$

$$= \frac{1}{2} m \dot{x}^2 + \frac{1}{2} m \dot{x}^2 + \frac{1}{2} m l^2 \dot{\theta}^2 (\sin^2 \theta + \cos^2 \theta)$$

$$- \frac{1}{2} m (2 \dot{x} l \sin \theta \dot{\theta})$$

$$= m \dot{x}^2 + \frac{1}{2} m l^2 \dot{\theta}^2 - m \dot{x} \dot{\theta} l \sin \theta \approx m \dot{x}^2 + \frac{1}{2} m l^2 \dot{\theta}^2 \text{ for small } \theta.$$



$V = \text{potential energy} = \text{potential energy of spring} + \text{potential energy of bob}$

$$= \frac{1}{2} k x^2 + m g l (1 - \cos \theta)$$

(Note: Potential energy of slider need not be considered if $x=0$ corresponds to static equilibrium position)

$$\text{Since } \cos \theta \approx 1 - \frac{1}{2} \theta^2, \quad V = \frac{1}{2} k x^2 + \frac{1}{2} m g l \theta^2$$

As there are no external forces, Lagrange's equations become

$$\frac{d}{dt} \left(\frac{\partial T}{\partial \dot{q}_j} \right) - \frac{\partial T}{\partial q_j} + \frac{\partial V}{\partial q_j} = 0; \quad j = 1, 2$$

Here $q_1 = x$ and $q_2 = \theta$

$$\frac{\partial T}{\partial x} = 0, \quad \frac{\partial T}{\partial \dot{x}} = 2 m \dot{x}, \quad \frac{d}{dt} \left(\frac{\partial T}{\partial \dot{x}} \right) = 2 m \ddot{x}, \quad \frac{\partial V}{\partial x} = k x$$

$$\frac{\partial T}{\partial \theta} = 0, \quad \frac{\partial T}{\partial \dot{\theta}} = m l^2 \dot{\theta}, \quad \frac{d}{dt} \left(\frac{\partial T}{\partial \dot{\theta}} \right) = m l^2 \ddot{\theta}, \quad \frac{\partial V}{\partial \theta} = m g l \theta$$

Lagrange's equations become

$$2 m \ddot{x} + k x = 0; \quad m l^2 \ddot{\theta} + m g l \theta = 0 \quad \text{or} \quad l \ddot{\theta} + g \theta = 0$$

5. P. 6.40

20 Pts

(1) With x_1 and x_2 as generalized coordinates:Since $x_1 = x - l_1 \theta$ and $x_2 = x + l_2 \theta$,

$$x = \left(\frac{x_1 l_2 + x_2 l_1}{l_1 + l_2} \right) \quad \text{and} \quad \theta = \left(\frac{x_2 - x_1}{l_1 + l_2} \right)$$

$$T = \frac{1}{2} m \dot{x}^2 + \frac{1}{2} J_0 \dot{\theta}^2 = \frac{1}{2} m \left(\frac{\dot{x}_1 l_2 + \dot{x}_2 l_1}{l_1 + l_2} \right)^2 + \frac{1}{2} J_0 \left(\frac{\dot{x}_2 - \dot{x}_1}{l_1 + l_2} \right)^2$$

$$\frac{\partial T}{\partial x_1} = 0, \quad \frac{\partial T}{\partial \dot{x}_1} = \frac{m}{2(l_1 + l_2)^2} (2 l_2^2 \dot{x}_1 + 2 l_1 l_2 \dot{x}_2) + \frac{J_0}{2(l_1 + l_2)^2} (2 \dot{x}_1 - 2 \dot{x}_2),$$

$$\frac{d}{dt} \left(\frac{\partial T}{\partial \dot{x}_1} \right) = \frac{m}{(l_1 + l_2)^2} (l_2^2 \ddot{x}_1 + l_1 l_2 \ddot{x}_2) + \frac{J_0}{(l_1 + l_2)^2} (\ddot{x}_1 - \ddot{x}_2)$$

$$\frac{\partial T}{\partial x_2} = 0, \quad \frac{\partial T}{\partial \dot{x}_2} = \frac{m}{2(l_1 + l_2)^2} (2 l_1^2 \dot{x}_2 + 2 l_1 l_2 \dot{x}_1) + \frac{J_0}{2(l_1 + l_2)^2} (2 \dot{x}_2 - 2 \dot{x}_1),$$

$$\frac{d}{dt} \left(\frac{\partial T}{\partial \dot{x}_2} \right) = \frac{m}{(l_1 + l_2)^2} (l_1^2 \ddot{x}_2 + l_1 l_2 \ddot{x}_1) + \frac{J_0}{(l_1 + l_2)^2} (\ddot{x}_2 - \ddot{x}_1)$$

$$V = \frac{1}{2} k_1 x_1^2 + \frac{1}{2} k_2 x_2^2$$

$$\frac{\partial V}{\partial x_1} = k_1 x_1, \quad \frac{\partial V}{\partial x_2} = k_2 x_2$$

Lagrange's equations, Eq. (6.41), give

$$\frac{m}{(l_1 + l_2)^2} (l_2^2 \ddot{x}_1 + l_1 l_2 \ddot{x}_2) + \frac{J_0}{(l_1 + l_2)^2} (\ddot{x}_1 - \ddot{x}_2) + k_1 x_1 = 0$$

$$\frac{m}{(l_1 + l_2)^2} (l_1^2 \ddot{x}_2 + l_1 l_2 \ddot{x}_1) + \frac{J_0}{(l_1 + l_2)^2} (\ddot{x}_2 - \ddot{x}_1) + k_2 x_2 = 0$$

$$\text{i.e.} \quad \ddot{x}_1 (m l_2^2 + J_0) + \ddot{x}_2 (m l_1 l_2 - J_0) + x_1 (l_1 + l_2)^2 k_1 = 0$$

$$\ddot{x}_1 (m l_1 l_2 - J_0) + \ddot{x}_2 (m l_1^2 + J_0) + x_2 (l_1 + l_2)^2 k_2 = 0$$

(2) With x and θ as generalized coordinates:

$$T = \frac{1}{2} m \dot{x}^2 + \frac{1}{2} J_0 \dot{\theta}^2$$

$$V = \frac{1}{2} k_1 (x - l_1 \theta)^2 + \frac{1}{2} k_2 (x + l_2 \theta)^2$$

$$\frac{\partial T}{\partial x} = 0, \quad \frac{\partial T}{\partial \dot{x}} = m \dot{x}, \quad \frac{d}{dt} \left(\frac{\partial T}{\partial \dot{x}} \right) = m \ddot{x}, \quad \frac{\partial V}{\partial x} = k_1 (x - l_1 \theta) + k_2 (x + l_2 \theta)$$

$$\frac{\partial T}{\partial \theta} = 0, \quad \frac{\partial T}{\partial \dot{\theta}} = J_0 \dot{\theta}, \quad \frac{d}{dt} \left(\frac{\partial T}{\partial \dot{\theta}} \right) = J_0 \ddot{\theta}, \quad \frac{\partial V}{\partial \theta} = -k_1 l_1 (x - l_1 \theta) + k_2 l_2 (x + l_2 \theta)$$

Lagrange's equations give

$$m \ddot{x} + k_1 (x - l_1 \theta) + k_2 (x + l_2 \theta) = 0$$

$$J_0 \ddot{\theta} - k_1 l_1 (x - l_1 \theta) + k_2 l_2 (x + l_2 \theta) = 0$$

$$[m] = \begin{bmatrix} m_1 & 0 & 0 \\ 0 & m_2 & 0 \\ 0 & 0 & m_3 \end{bmatrix} = \begin{bmatrix} m & 0 & 0 \\ 0 & 2m & 0 \\ 0 & 0 & 3m \end{bmatrix}, \quad [k] = \begin{bmatrix} k_1 & -k_1 & 0 \\ -k_1 & (k_1+k_2) & -k_2 \\ 0 & -k_2 & k_2 \end{bmatrix} = \begin{bmatrix} k & -k & 0 \\ -k & 2k & -k \\ 0 & -k & k \end{bmatrix}$$

Equations of motion for harmonic motion:

$$\begin{bmatrix} (-m\omega^2 + k) & -k & 0 \\ -k & (-2m\omega^2 + 2k) & -k \\ 0 & -k & (-3m\omega^2 + k) \end{bmatrix} \begin{Bmatrix} X_1 \\ X_2 \\ X_3 \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \\ 0 \end{Bmatrix} \quad (E.1)$$

Defining $\alpha = \frac{m\omega^2}{k}$, (E.1) can be rewritten as

$$\begin{bmatrix} (-\alpha + 1) & -1 & 0 \\ -1 & (-2\alpha + 2) & -1 \\ 0 & -1 & (-3\alpha + 1) \end{bmatrix} \begin{Bmatrix} X_1 \\ X_2 \\ X_3 \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \\ 0 \end{Bmatrix} \quad (E.2)$$

Frequency equation is

$$2\alpha(3\alpha^2 - 7\alpha + 3) = 0$$

Roots are

$$\begin{aligned} \alpha_1 &= 0 & ; & \quad \omega_1 = 0 \\ \alpha_2 &= 0.565741 & ; & \quad \omega_2 = 0.752158 \sqrt{k/m} \\ \alpha_3 &= 1.767592 & ; & \quad \omega_3 = 1.329508 \sqrt{k/m} \end{aligned}$$

Eqs. (E.2) give $X_2^{(j)} = (-\alpha_j + 1) X_1^{(j)}, \quad X_3^{(j)} = \left(\frac{1}{-3\alpha_j + 1} \right) X_2^{(j)}$

$$\therefore \vec{X}^{(j)} = \begin{Bmatrix} 1.0 \\ (-\alpha_j + 1) \\ \left(\frac{-\alpha_j + 1}{-3\alpha_j + 1} \right) \end{Bmatrix} X_1^{(j)}$$

Hence

$$\vec{X}^{(1)} = \begin{Bmatrix} 1 \\ 1 \\ 1 \end{Bmatrix} X_1^{(1)}, \quad \vec{X}^{(2)} = \begin{Bmatrix} 1 \\ 0.434259 \\ -0.622841 \end{Bmatrix} X_1^{(2)}, \quad \vec{X}^{(3)} = \begin{Bmatrix} 1 \\ -0.767592 \\ 0.178395 \end{Bmatrix} X_1^{(3)}$$

6.77

From solution of Problem 6.57, the natural frequencies and mode shapes are given by:

$$\omega_1 = 0.482087 \sqrt{\frac{k}{m}} ; \omega_2 = \sqrt{\frac{k}{m}} ; \omega_3 = 1.197605 \sqrt{\frac{k}{m}}$$

$$\vec{X}^{(1)} = \begin{Bmatrix} 1.0 \\ 1.535184 \\ 1.0 \end{Bmatrix} ; \vec{X}^{(2)} = \begin{Bmatrix} 1 \\ 0 \\ -1 \end{Bmatrix} ; \vec{X}^{(3)} = \begin{Bmatrix} 1.0 \\ 0.868516 \\ -1 \end{Bmatrix}$$

Initial conditions:

$$x_1(0) = x_{10}, x_2(0) = 0, \dot{x}_3(0) = 0, \dot{x}_i(0) = 0 ; i = 1, 2, 3$$

Equations (6.98) and (6.99) yield:

$$A_1 \cos \phi_1 + A_2 \cos \phi_2 + A_3 \cos \phi_3 = x_{10} \quad (1)$$

$$1.5352 A_1 \cos \phi_1 + 0.8685 A_3 \cos \phi_3 = 0 \quad (2)$$

$$A_1 \cos \phi_1 - A_2 \cos \phi_2 - A_3 \cos \phi_3 = 0 \quad (3)$$

$$-0.4821 \sqrt{\frac{k}{m}} A_1 \sin \phi_1 - \sqrt{\frac{k}{m}} A_2 \sin \phi_2 - 1.1976 \sqrt{\frac{k}{m}} A_3 \sin \phi_3 = 0 \quad (4)$$

$$-0.4821 \sqrt{\frac{k}{m}} (1.5352) A_1 \sin \phi_1 - 1.1976 \sqrt{\frac{k}{m}} (0.8685) A_3 \sin \phi_3 = 0 \quad (5)$$

$$-0.4821 \sqrt{\frac{k}{m}} (1.0) A_1 \sin \phi_1 - \sqrt{\frac{k}{m}} (-1) A_2 \sin \phi_2 - 1.1976 \sqrt{\frac{k}{m}} (-1) A_3 \sin \phi_3 = 0 \quad (6)$$

Solution of Eqs. (4) to (6):

$$\phi_i = 0 ; i = 1, 2, 3$$

Solution of Eqs. (1) to (3) gives:

$$A_1 = 0.5 x_{10} ; A_2 = 0.3838 x_{10} ; A_3 = -0.8838 x_{10}$$

Thus the free vibration solution of the system is given by:

$$x_1(t) = x_{10} (0.5 \cos 0.4821 \sqrt{\frac{k}{m}} t + 0.3838 \cos \sqrt{\frac{k}{m}} t - 0.8838 \cos 1.1976 \sqrt{\frac{k}{m}} t)$$

$$x_2(t) = x_{10} \left\{ 0.7676 \cos 0.4821 \sqrt{\frac{k}{m}} t - 0.7676 \cos 1.1976 \sqrt{\frac{k}{m}} t \right\}$$

$$x_3(t) = x_{10} (0.5 \cos 0.4821 \sqrt{\frac{k}{m}} t - 0.3838 \cos \sqrt{\frac{k}{m}} t + 0.8838 \cos 1.1976 \sqrt{\frac{k}{m}} t)$$

Equations of motion are

6.89
$$\begin{bmatrix} 2m & 0 & 0 \\ 0 & m & 0 \\ 0 & 0 & m \end{bmatrix} \begin{Bmatrix} \ddot{x}_1 \\ \ddot{x}_2 \\ \ddot{x}_3 \end{Bmatrix} + \begin{bmatrix} 3k & -k & -k \\ -k & k & 0 \\ -k & 0 & k \end{bmatrix} \begin{Bmatrix} x_1 \\ x_2 \\ x_3 \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \\ F(t) \end{Bmatrix}$$

where $m = 1 \text{ kg}$, $k = 1000 \text{ N/m}$, $F(t) = 5 \sin 10t \text{ N}$.

Eigenvalue analysis:

Frequency equation is

$$\left| -\omega^2 m \begin{bmatrix} 2 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} + k \begin{bmatrix} 3 & -1 & -1 \\ -1 & 1 & 0 \\ -1 & 0 & 1 \end{bmatrix} \right| = 0$$

$$\text{or } \begin{vmatrix} -2\lambda + 3 & -1 & -1 \\ -1 & -\lambda + 1 & 0 \\ -1 & 0 & -\lambda + 1 \end{vmatrix} = 0 \quad \text{where } \lambda = \frac{\omega^2 m}{k}$$

$$\text{or } 2\lambda^3 - 7\lambda^2 + 6\lambda - 1 = 0$$

Roots are

$$\lambda_1 = 0.219220, \quad \lambda_2 = 1.0, \quad \lambda_3 = 2.28078$$

$$\omega_1 = 0.4682094 \sqrt{\frac{k}{m}}, \quad \omega_2 = 1.0 \sqrt{\frac{k}{m}}, \quad \omega_3 = 1.5102251 \sqrt{\frac{k}{m}}$$

Mode shapes are given by

$$\begin{bmatrix} -2\lambda_i + 3 & -1 & -1 \\ -1 & -\lambda_i + 1 & 0 \\ -1 & 0 & -\lambda_i + 1 \end{bmatrix} \begin{Bmatrix} x_1^{(i)} \\ x_2^{(i)} \\ x_3^{(i)} \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \\ 0 \end{Bmatrix}$$

$$\text{or } x_1^{(i)} = (-\lambda_i + 1) x_2^{(i)}, \quad x_3^{(i)} = (-2\lambda_i + 3) x_1^{(i)} - x_2^{(i)}$$

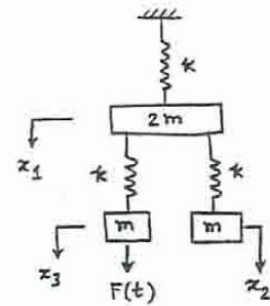
$$\vec{X}^{(i)} = \begin{Bmatrix} -\lambda_i + 1 \\ 1 \\ (-2\lambda_i + 3)(-\lambda_i + 1) - 1 \end{Bmatrix} x_2^{(i)}$$

$$\vec{X}^{(1)} = \begin{Bmatrix} 0.78078 \\ 1 \\ 1 \end{Bmatrix} x_2^{(1)}, \quad \vec{X}^{(2)} = \begin{Bmatrix} 0 \\ 1 \\ -1 \end{Bmatrix} x_2^{(2)}, \quad \vec{X}^{(3)} = \begin{Bmatrix} -1.28078 \\ 1 \\ 1 \end{Bmatrix} x_2^{(3)}$$

Normalization of mode shapes with respect to $[m]$:

$$\vec{X}^{(1)T} [m] \vec{X}^{(1)} = (0.78078 \quad 1 \quad 1) \begin{Bmatrix} 1.56156 \\ 1 \\ 1 \end{Bmatrix} (x_2^{(1)})^2 = 3.21923 (x_2^{(1)})^2 = 1$$

$$x_2^{(1)} = 0.55734; \quad \vec{X}^{(1)} = \begin{Bmatrix} 0.43516 \\ 0.55734 \\ 0.55734 \end{Bmatrix}$$



$$\vec{x}^{(2)T} [m] \vec{x}^{(2)} = (0 \quad 1 \quad -1) \begin{Bmatrix} 0 \\ 1 \\ -1 \end{Bmatrix} (x_2^{(2)})^2 = 2 (x_2^{(2)})^2 = 1$$

$$x_2^{(2)} = 0.70711; \quad \vec{x}^{(2)} = \begin{Bmatrix} 0 \\ 0.70711 \\ -0.70711 \end{Bmatrix}$$

$$\vec{x}^{(3)T} [m] \vec{x}^{(3)} = (-1.28078 \quad 1 \quad 1) \begin{Bmatrix} -2.56156 \\ 1 \\ 1 \end{Bmatrix} (x_2^{(3)})^2 = 5.28079 (x_2^{(3)})^2 = 1$$

$$x_2^{(3)} = 0.43516; \quad \vec{x}^{(3)} = \begin{Bmatrix} -0.55734 \\ 0.43516 \\ 0.43516 \end{Bmatrix}$$

$\vec{Q} = [X]^T \vec{F}(t)$ = vector of generalized forces

$$= \begin{bmatrix} 0.43516 & 0.55734 & 0.55734 \\ 0 & 0.70711 & -0.70711 \\ -0.55734 & 0.43516 & 0.43516 \end{bmatrix} \begin{Bmatrix} 0 \\ 0 \\ F_0 \sin \omega t \end{Bmatrix} = \begin{Bmatrix} 2.7867 \\ -3.5356 \\ 2.1758 \end{Bmatrix} \sin 10t$$

Uncoupled equations of motion are

$$\left. \begin{aligned} \ddot{v}_1 + 219.22 v_1 &= 2.7867 \sin 10t \\ \ddot{v}_2 + 1000.00 v_2 &= -3.5356 \sin 10t \\ \ddot{v}_3 + 2280.78 v_3 &= 2.1758 \sin 10t \end{aligned} \right\} \quad (E.1)$$

Particular solutions of (E.1) are

$$v_1(t) = \left(\frac{2.7867}{219.22 - 100} \right) \sin 10t = 0.0233744 \sin 10t$$

$$v_2(t) = \left(\frac{-3.5356}{1000 - 100} \right) \sin 10t = -0.0039284 \sin 10t$$

$$v_3(t) = \left(\frac{2.1758}{2280.78 - 100} \right) \sin 10t = 0.0009977 \sin 10t$$

Since $\vec{x} = [X] \vec{v} = \begin{bmatrix} 0.43516 & 0 & -0.55734 \\ 0.55734 & 0.70711 & 0.43516 \\ 0.55734 & -0.70711 & 0.43516 \end{bmatrix} \vec{v}$, we get

$$x_1(t) = 0.0096155 \sin 10t \quad m$$

$$x_2(t) = 0.0095333 \sin 10t \quad m$$

$$x_3(t) = 0.0162395 \sin 10t \quad m.$$